

# $H \rightarrow \tau^+ \tau^-$ branching ratio study at $\sqrt{s} = 250$ GeV at the ILC with the ILD detector

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## Abstract

We evaluated the measurement accuracy of the branching ratio of  $H \rightarrow \tau^+ \tau^-$  mode at  $\sqrt{s} = 250$  GeV at the ILC with the ILD detector. We assumed the Higgs mass  $M_H = 120$  GeV, branching ratio  $\text{Br}(H \rightarrow \tau^+ \tau^-) = 8.0$  %, beam polarization  $P(e^-, e^+) = (-0.8, +0.3)$ , and integrated luminosity  $\int L dt = 250 \text{ fb}^{-1}$ . We used the LOI samples as the Monte-Carlo samples. The evaluation was performed by the ILD full detector simulation. All Standard Model backgrounds were included in this study. We obtained the accuracy  $\Delta(\sigma \cdot \text{Br})/(\sigma \cdot \text{Br}) = 3.5$  %. The scaled result to  $M_H = 125$  GeV is calculated to be 4.2 %.

## 1 Introduction

A new Higgs-like particle was discovered by the ATLAS and the CMS experiments [1, 2]. One of the next important themes for particle physics is the investigation of that new particle, especially the mass generation mechanism.

One of the important properties of Higgs boson is its branching ratio. In the Standard Model (SM) of particle physics, the Yukawa coupling constant of matter fermions with the Higgs boson is proportional to the fermion mass. Besides, if there is new physics, the coupling constant may deviate from the SM prediction. Therefore, the branching ratio is a probe for new physics.

In this note, we focus on the branching ratio of  $H \rightarrow \tau^+ \tau^-$  mode. We estimate the measurement accuracy of the  $H \rightarrow \tau^+ \tau^-$  branching ratio at  $\sqrt{s} = 250$  GeV with the ILD full detector simulation.

## 2 Signal and Background

The main Higgs production process at  $\sqrt{s} = 250$  GeV is the Higgs-strahlung process ( $e^+ e^- \rightarrow ZH$ ). There are three types of signal depending on the decay of the  $Z$  boson, as shown in Figure 1. In this note, we concentrate on (A)  $Z \rightarrow l^+ l^-$  mode and (B)  $Z \rightarrow q \bar{q}$  mode.

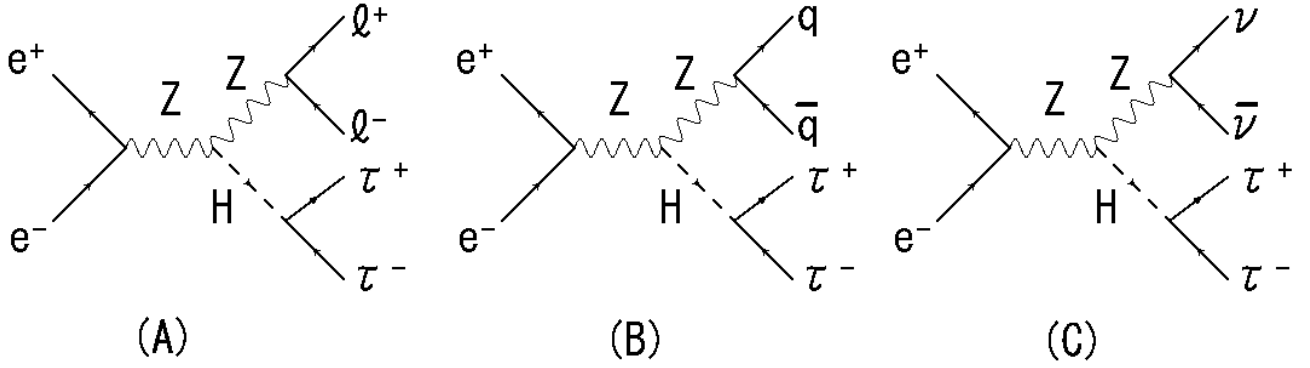


Figure 1: The diagrams of signal processes. (A):  $Z \rightarrow l^+l^-$  mode, (B):  $Z \rightarrow q\bar{q}$  mode, (C):  $Z \rightarrow \nu\bar{\nu}$  mode.

The  $Z \rightarrow \nu\bar{\nu}$  mode has been found to contribute negligibly to the overall precision which is dominated by the  $Z \rightarrow q\bar{q}$  mode. However, at higher center-of-mass energies, the  $e^+e^- \rightarrow \nu\bar{\nu}H$  mode is expected to contribute substantially due to the increase in the cross section of  $WW$  fusion process.

## 2.1 $Z \rightarrow l^+l^-$ mode

In this mode, we only considered  $Z \rightarrow e^+e^-$  mode and  $Z \rightarrow \mu^+\mu^-$  mode as the signal process. The signal cross section of this mode is 1.9 fb. The dominant background processes are the four leptons processes ( $e^+e^- \rightarrow eeee$ ,  $ee\mu\mu$ ,  $ee\tau\tau$ ,  $\mu\mu\mu\mu$ ,  $\mu\mu\tau\tau$ , and  $\tau\tau\tau\tau$ ). An example diagram is shown in Figure 2-(A). Other background processes are  $e^+e^- \rightarrow ZH$  reactions where the Higgs boson does not decay to tau pairs.

## 2.2 $Z \rightarrow q\bar{q}$ mode

The signal cross section of this mode is 19.8 fb. The possible background processes for this mode are  $qqqq$ ,  $qqll$ , and  $qq\nu\nu$ , which come from  $e^+e^- \rightarrow W^+W^-$  or  $e^+e^- \rightarrow ZZ$  reactions. An example diagram is shown in Figure 2-(B). Other possible backgrounds are  $e^+e^- \rightarrow ZH$  with  $Z \rightarrow \tau^+\tau^-$  and  $H \rightarrow q\bar{q}$ . These processes have the same final state to the signal.

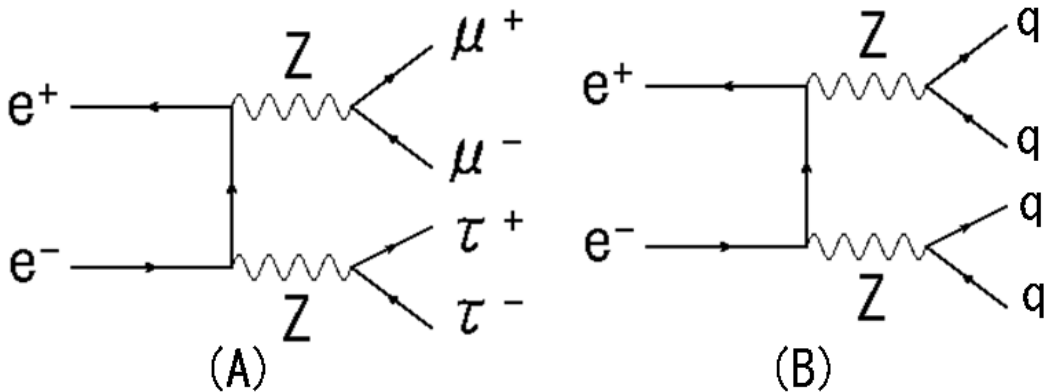


Figure 2: Example diagrams of possible background. (A):  $\mu\mu\tau\tau$  background for  $Z \rightarrow l^+l^-$  mode, (B):  $qqqq$  background for  $Z \rightarrow q\bar{q}$  mode.

### 3 Simulation Conditions

We performed the detector simulation with Mokka [3], a Geant4-based [4] full simulation, with the ILD\_00 detector model. TAUOLA [5] was used for the tau decay simulation. The ILD\_00 detector model consists of vertex detector, time projection chamber, electromagnetic calorimeter (ECAL), hadronic calorimeter (HCAL), and yoke.

We used the signal and background samples which were generated in the context of the Letter of Intent [6]. The assumed center-of-mass energy is 250 GeV. The effects of beamstrahlung and initial state radiation are included. All Monte-Carlo sample information (process ID, process, polarization, cross section, number of events, and luminosity) are summarized in Tables 6 (page 9) and 7 (page 10). We assumed the Higgs mass  $M_H = 120$  GeV, branching ratio  $\text{Br}(H \rightarrow \tau^+\tau^-) = 8.0\%$  as assumed by PYTHIA [8], integrated luminosity  $\int L dt = 250 \text{ fb}^{-1}$ , and beam polarization  $P(e^+, e^-) = (+0.3, -0.8)$ . We also rescale the final result to the case of  $M_H = 125$  GeV and the  $H \rightarrow \tau^+\tau^-$  branching ratio which includes the NNLO corrections [9].

### 4 Event Reconstruction and Event Selection

#### 4.1 $Z \rightarrow l^+l^-$ mode

In this mode, we take the strategy of reconstructing the  $Z$  boson first, followed by the reconstruction of the tau pairs from the Higgs decay.

We applied lepton identification at first for dividing  $Z \rightarrow e^+e^-$  events and  $Z \rightarrow \mu^+\mu^-$  events by using the information of energy deposit in the calorimeter ( $E_{\text{ECAL}}$  and  $E_{\text{HCAL}}$ , where  $E_{\text{ECAL}}$  is the energy deposit in ECAL,  $E_{\text{HCAL}}$  is the energy deposit in HCAL, respectively) and track momentum ( $P_{\text{track}}$ ). Figures 3 - 6 are the plots of  $E_{\text{ECAL}}/(E_{\text{ECAL}} + E_{\text{HCAL}})$  and  $(E_{\text{ECAL}} + E_{\text{HCAL}})/P_{\text{track}}$ .

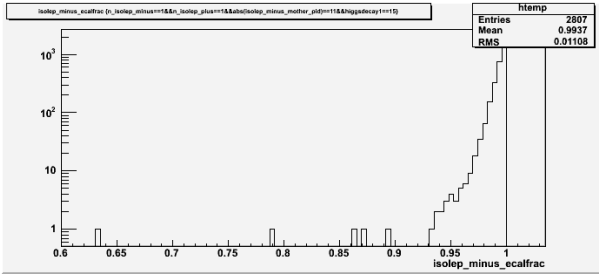


Figure 3: The plot of  $E_{\text{ECAL}}/(E_{\text{ECAL}} + E_{\text{HCAL}})$  for the  $e$  in  $eeH$  samples.

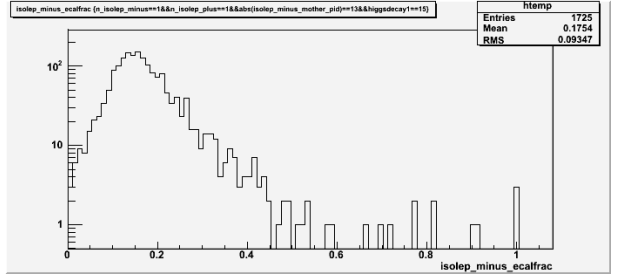


Figure 4: The plot of  $E_{\text{ECAL}}/(E_{\text{ECAL}} + E_{\text{HCAL}})$  for the  $\mu$  in  $\mu\mu H$  samples.

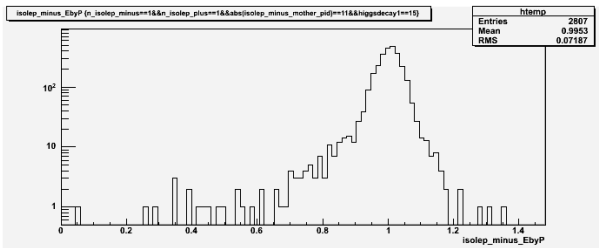


Figure 5: The plot of  $(E_{\text{ECAL}} + E_{\text{HCAL}})/P_{\text{track}}$  for the  $e$  in  $eeH$  samples.

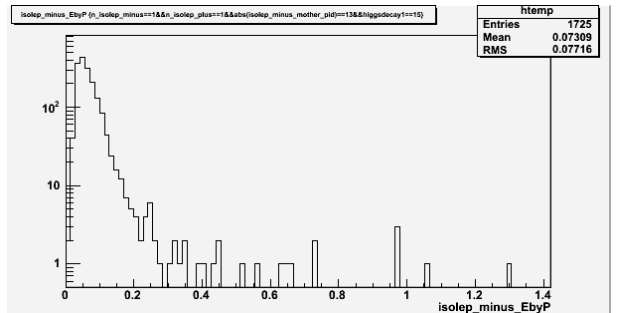


Figure 6: The plot of  $(E_{\text{ECAL}} + E_{\text{HCAL}})/P_{\text{track}}$  for the  $\mu$  in  $\mu\mu H$  samples.

From these plots, we define the criteria for lepton identification. The criteria for electron identification ( $e$ -ID) are:  $E_{\text{ECAL}}/(E_{\text{ECAL}} + E_{\text{HCAL}}) > 0.92$  and  $(E_{\text{ECAL}} + E_{\text{HCAL}})/P_{\text{track}} > 0.5$ . The criteria for muon identification ( $\mu$ -ID) are:  $E_{\text{ECAL}}/(E_{\text{ECAL}} + E_{\text{HCAL}}) < 0.6$  and  $(E_{\text{ECAL}} + E_{\text{HCAL}})/P_{\text{track}} < 0.5$ .

After the lepton identification, we applied selections to remove secondary leptons from tau decays. The strategy of this selection is to remove tracks which do not come from the interaction point (IP) by using the track energy  $E_{\text{track}}$  and impact parameter in the transverse direction  $d_0$  and longitudinal direction  $z_0$  with respect to the beam axis. Figures 7 - 12 show the  $|d_0/\sigma(d_0)|$ ,  $|z_0/\sigma(z_0)|$ , and  $E_{\text{track}}$  plots which through the lepton identification. We defined the tau rejection cut for the objects through the  $e$ -ID:  $|d_0/\sigma(d_0)| < 50$ ,  $|z_0/\sigma(z_0)| < 5$ , and  $E_{\text{track}} > 10$  GeV, and for the objects through the  $\mu$ -ID:  $|d_0/\sigma(d_0)| < 3$ ,  $|z_0/\sigma(z_0)| < 7$ , and  $E_{\text{track}} > 20$  GeV.

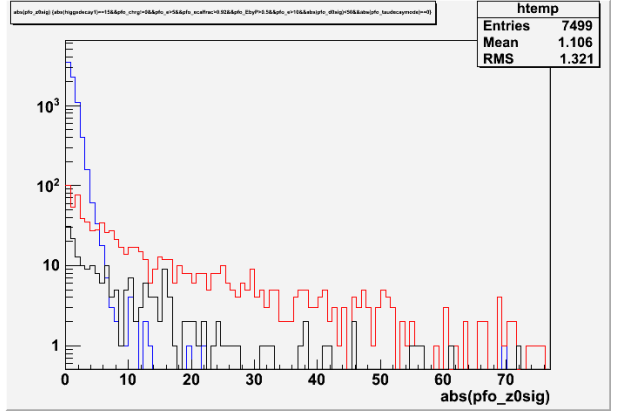
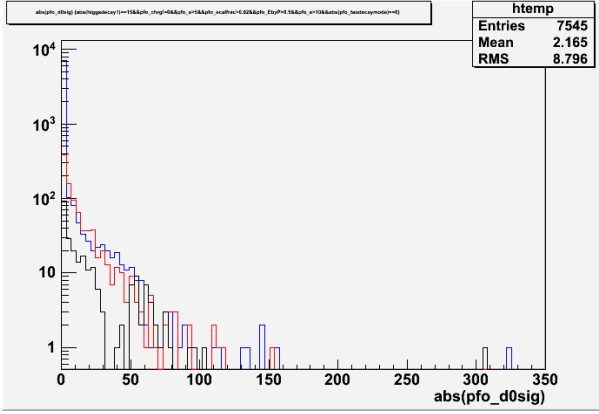


Figure 7: The plot of  $|d_0/\sigma(d_0)|$  of  $e$  of  $eeH$  process. Blue, red, and black histograms show the  $e$  from  $Z \rightarrow e^+e^-$ , the  $e$  from  $\tau \rightarrow e\nu\nu$ , and the hadrons from  $\tau$  decay, respectively.

Figure 8: The plot of  $|z_0/\sigma(z_0)|$  of  $e$  of  $eeH$  process. Blue, red, and black histograms show the  $e$  from  $Z \rightarrow e^+e^-$ , the  $e$  from  $\tau \rightarrow e\nu\nu$ , and the hadrons from  $\tau$  decay, respectively.

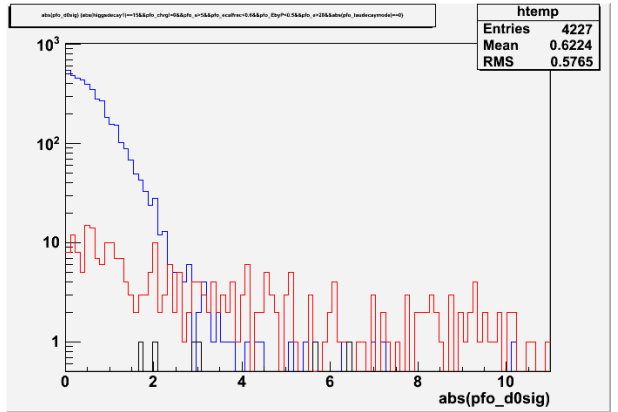
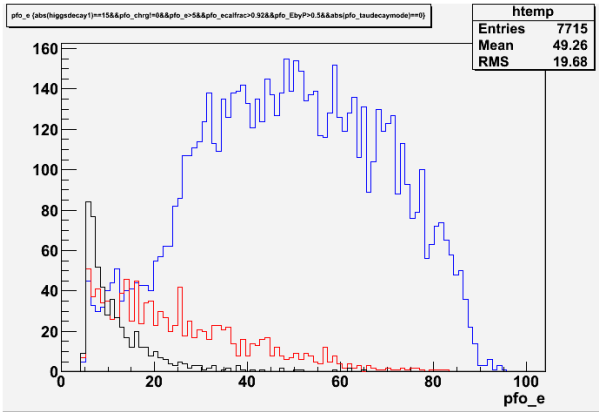


Figure 9: The plot of  $E_{\text{track}}$  of  $e$  of  $eeH$  process. Blue, red, and black histograms show the  $e$  from  $Z \rightarrow e^+e^-$ , the  $e$  from  $\tau \rightarrow e\nu\nu$ , and the hadrons from  $\tau$  decay, respectively.

Figure 10: The plot of  $|d_0/\sigma(d_0)|$  of  $\mu$  of  $\mu\mu H$  process. Blue, red, and black histograms show the  $\mu$  from  $Z \rightarrow \mu^+\mu^-$ , the  $\mu$  from  $\tau \rightarrow \mu\nu\nu$ , and the hadrons from  $\tau$  decay, respectively.

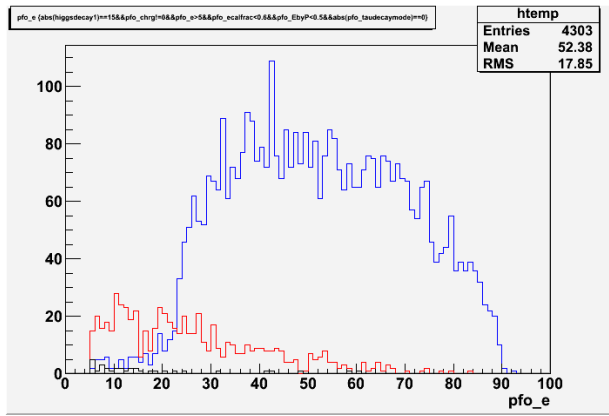
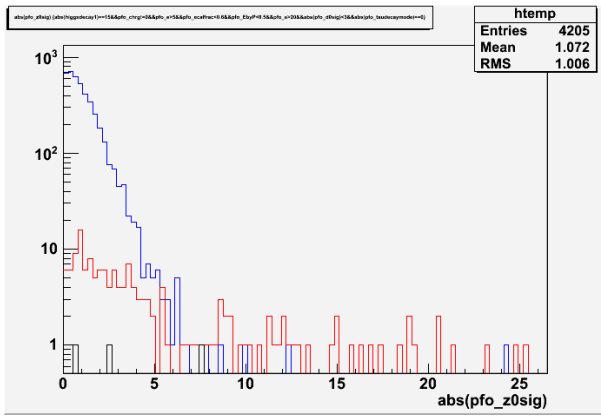


Figure 11: The plot of  $|z_0/\sigma(z_0)|$  of  $\mu$  of  $\mu\mu H$  process. Blue, red, and black histograms show the  $\mu$  from  $Z \rightarrow \mu^+\mu^-$ , the  $\mu$  from  $\tau \rightarrow \mu\nu\nu$ , and the hadrons from  $\tau$  decay, respectively. Figure 12: The plot of  $E_{\text{track}}$  of  $\mu$  of  $\mu\mu H$  process. Blue, red, and black histograms show the  $\mu$  from  $Z \rightarrow \mu^+\mu^-$ , the  $\mu$  from  $\tau \rightarrow \mu\nu\nu$ , and the hadrons from  $\tau$  decay, respectively.

We applied the energy recovery procedure to correct for bremsstrahlung and final state radiation. In order to reconstruct the original  $Z$  boson, we have to use both the charged particles and the radiated photons. To achieve this, we defined the cone as shown in Figure 13. The four-momenta of the neutral particles in the cone were combined with that of the lepton candidate. We defined the half-opening angle of the cone with  $\cos\theta_{\text{cone}} = 0.999$  and applied the recovery procedure to the lepton candidates. The results are shown in Figures 14 and 15.

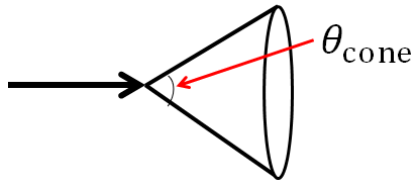


Figure 13: The definition of the cone. Black arrow shows the lepton candidate.  $\theta_{\text{cone}}$  is the angle of the cone.

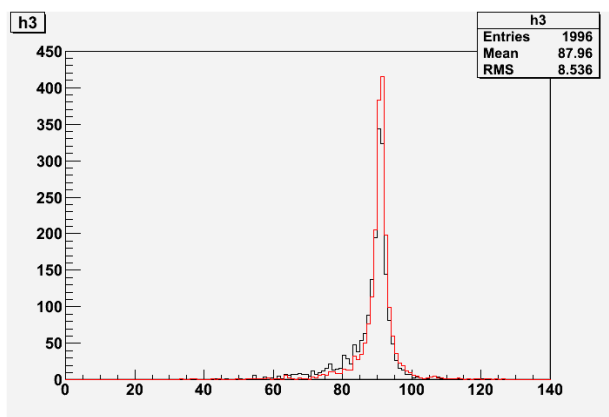
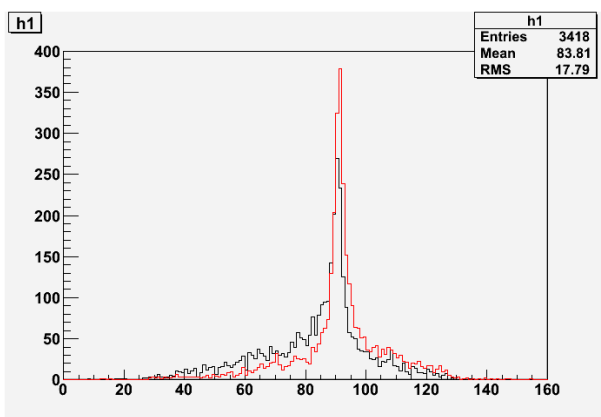


Figure 14: The results of recovery for  $Z \rightarrow e^+e^-$  mode. The horizontal axis shows the  $M_Z$ . Black and red histograms show the results of without recovery and with recovery ( $\cos\theta_{\text{cone}} = 0.999$ ), respectively. Figure 15: The results of recovery for  $Z \rightarrow \mu^+\mu^-$  mode. The horizontal axis shows the  $M_Z$ . Black and red histograms show the results of without recovery and with recovery ( $\cos\theta_{\text{cone}} = 0.999$ ), respectively.

After that, we applied the tau finder to the remaining objects to reconstruct tau leptons. First of all, the objects which already used at  $Z$  boson reconstruction were rejected from tau reconstruction analysis. Then we search the highest energy track from the remaining objects, and combine the neighboring particles (which satisfies the angle with respect to the highest energy track less than 1.0 radian) with the combined mass less than 2 GeV. We regarded the combined object as a tau candidate. Then repeat these processes until there are no charged particles.

After finishing the event reconstruction, we applied the cuts for selecting signal, rejecting background. Before optimizing the cuts, we applied the preselection as follows for  $Z \rightarrow e^+e^-$  mode: number of  $e^+$  and  $e^- = 1$ , number of  $\tau^+$  and  $\tau^- = 1$ , and for  $Z \rightarrow \mu^+\mu^-$  mode: number of  $\mu^+$  and  $\mu^- = 1$ , number of  $\tau^+$  and  $\tau^- = 1$ .

We applied the following cuts for  $Z \rightarrow e^+e^-$  mode: number of tracks  $\leq 8$ ,  $115 \text{ GeV} < E_{\text{vis}} < 230 \text{ GeV}$ ,  $|\cos \theta_{\text{miss}}| < 0.99$ ,  $81 \text{ GeV} < M_Z < 113 \text{ GeV}$ ,  $\cos \theta_{e^-} < 0.92$ ,  $\cos \theta_{e^+} > -0.92$ ,  $E_{e^-} < 90 \text{ GeV}$ ,  $E_{e^+} < 90 \text{ GeV}$ ,  $\cos \theta_{\tau^+\tau^-} < -0.45$ ,  $\cos \theta_{\tau^-} < 0.92$ ,  $\cos \theta_{\tau^+} > -0.92$ , and  $116 \text{ GeV} < M_{\text{recoil}} < 142 \text{ GeV}$ , where  $E_{\text{vis}}$  is the visible energy,  $\theta_{\text{miss}}$  is the missing momentum angle with respect to beam axis,  $\theta_{e^-(e^+)}$  is the  $e^-(e^+)$  angle with respect to beam axis,  $E_{e^-(e^+)}$  is the  $e^-(e^+)$  energy,  $\theta_{\tau^+\tau^-}$  is the angle between  $\tau^+$  and  $\tau^-$ ,  $\theta_{\tau^-(\tau^+)}$  is the  $\tau^-(\tau^+)$  angle with respect to beam axis, and  $M_{\text{recoil}}$  is the recoil mass, respectively. The histograms of all cut variables are shown in Figures 17 - 28 (page 11 - 12). Table 1 shows the cut statistics of this mode. After the cuts, the  $Z \rightarrow e^+e^-$  signal events of 108.9 and background events of 76.0 remained. The statistical significance was calculated to be  $S/\sqrt{S+B} = 108.9/\sqrt{108.9+76.0} = 8.0\sigma$ .

We applied the following cuts for  $Z \rightarrow \mu^+\mu^-$  mode: number of tracks  $\leq 8$ ,  $115 \text{ GeV} < E_{\text{vis}} < 235 \text{ GeV}$ ,  $|\cos \theta_{\text{miss}}| < 0.98$ ,  $72 \text{ GeV} < M_Z < 107 \text{ GeV}$ ,  $E_{e^-} < 90 \text{ GeV}$ ,  $E_{e^+} < 90 \text{ GeV}$ ,  $\cos \theta_{\tau^+\tau^-} < -0.5$ , and  $118 \text{ GeV} < M_{\text{recoil}} < 143 \text{ GeV}$ . The histograms of all cut variables are shown in Figures 29 - 36 (page 13 - 14). Table 2 shows the cut statistics of this mode. For the  $Z \rightarrow \mu^+\mu^-$  mode case, 131.2 signal events and 91.2 background events were remained. The significance was  $S/\sqrt{S+B} = 131.2/\sqrt{131.2+91.2} = 8.8\sigma$ .

Table 1: The cut statistics of  $Z \rightarrow e^+e^-$  mode.

	$eeH$ $H \rightarrow \tau\tau$	$\mu\mu H$ $H \rightarrow \tau\tau$	$\tau\tau H$ $H \rightarrow \tau\tau$	$ZH$ with no $\tau$	$ee\tau\tau$	other 4 leptons	other SM bkg	signi.
No cut	228.3	211.1	214.6	7325	$2.388 \times 10^5$	$5.238 \times 10^5$	$1.492 \times 10^{10}$	0.0019
preselection	171.3	0.155	1.532	47.05	$1.338 \times 10^4$	$3.215 \times 10^4$	$1.023 \times 10^7$	0.053
# of tracks	169.4	0.155	1.532	41.56	$1.316 \times 10^4$	$3.205 \times 10^4$	$1.009 \times 10^7$	0.053
$E_{\text{vis}}$	162.3	0.155	0.912	38.36	$1.068 \times 10^4$	$1.039 \times 10^4$	$4.761 \times 10^6$	0.074
$\cos \theta_{\text{miss}}$	160.6	0.155	0.912	38.03	8719	1906	$5.155 \times 10^5$	0.22
$M_Z$	148.0	0	0.017	29.09	2408	501.2	$1.299 \times 10^4$	1.2
$\cos \theta_{e^-(e^+)}$	133.9	0	0.009	25.40	1067	101.5	729.7	3.0
$E_{e^-(e^+)}$	133.0	0	0.009	24.93	690.3	78.70	629.7	3.4
$\cos \theta_{\tau^+\tau^-}$	130.8	0	0	3.536	254.9	30.70	155.4	5.5
$\cos \theta_{\tau^-(\tau^+)}$	123.4	0	0	3.074	212.1	9.161	3.817	6.6
$M_{\text{recoil}}$	108.9	0	0	2.474	72.35	1.134	0.034	8.0

Table 2: The cut statistics of  $Z \rightarrow \mu^+\mu^-$  mode.

	$\mu\mu H$ $H \rightarrow \tau\tau$	$eeH$ $H \rightarrow \tau\tau$	$\tau\tau H$ $H \rightarrow \tau\tau$	$ZH$ with no $\tau$	$\mu\mu\tau\tau$	other 4 leptons	other SM bkg	signi.
No cut	211.1	228.3	214.6	7325	3513	$7.591 \times 10^6$	$1.492 \times 10^{10}$	0.0017
preselection	168.5	0	0.155	43.01	1698	7546	7732	1.3
# of tracks	167.4	0	0.155	39.65	1684	7537	7400	1.3
$E_{\text{vis}}$	162.9	0	0.155	37.40	1586	2285	3713	1.9
$\cos \theta_{\text{miss}}$	158.6	0	0.155	36.51	1386	227.5	55.48	3.7
$M_Z$	153.2	0	0	32.84	1038	55.28	42.54	4.2
$E_{e^-(e^+)}$	153.2	0	0	32.70	738.6	42.41	36.72	4.8
$\cos \theta_{\tau^+\tau^-}$	146.3	0	0	3.638	259.4	20.19	0.756	7.1
$M_{\text{recoil}}$	131.2	0	0	2.875	82.36	5.311	0.301	8.8



tively. The histograms of all cut variables are shown in Figures 37 - 49 (page 14 - 16). Table 3 shows the cut statistics of this mode. After the cuts, the signal events and background events were remained 1026 and 554.4. The statistical significance of  $Z \rightarrow q\bar{q}$  mode is calculated to be  $S/\sqrt{S+B} = 1026/\sqrt{1026+554.4} = 25.8\sigma$ .

Table 3: The cut statistics of  $Z \rightarrow q\bar{q}$  mode.

	$qqH$ $H \rightarrow \tau\tau$	$ZH$ with no $\tau$	$lH$	$\tau\tau H$	$qqqq$	$qql$	$qq\tau\tau$	$qql\nu$	$qq\tau\nu$	other SM bkg	signi.
No cut	4233	$4.829 \times 10^4$	5377	2596	$4.038 \times 10^6$	$3.563 \times 10^5$	$4.169 \times 10^4$	$2.788 \times 10^6$	$1.326 \times 10^6$	$1.494 \times 10^{10}$	0.035
preselection	1647	578.8	2761	765.4	$1.230 \times 10^4$	$6.378 \times 10^4$	$1.161 \times 10^4$	$1.249 \times 10^5$	$4.948 \times 10^4$	$2.570 \times 10^7$	0.32
# of tracks	1644	549.8	2680	765.4	$1.230 \times 10^4$	$6.059 \times 10^4$	$1.146 \times 10^4$	$1.214 \times 10^5$	$4.806 \times 10^4$	$5.190 \times 10^5$	1.9
$E_{vis}$	1607	492.3	1015	744.2	4443	$2.106 \times 10^4$	$1.107 \times 10^4$	$1.192 \times 10^5$	$4.693 \times 10^4$	$2.383 \times 10^5$	2.4
$\cos\theta_{miss}$	1572	474.7	860.5	725.1	2127	8315	$1.021 \times 10^4$	$1.171 \times 10^5$	$4.415 \times 10^4$	5939	3.6
$M_Z$	1440	376.1	791.3	682.8	778.6	4987	8674	8189	3288	997.3	8.3
$E_Z$	1429	352.0	782.7	528.7	505.0	4797	7857	7703	3061	609.9	8.6
$\cos\theta_{\tau^+\tau^-}$	1386	46.28	442.2	255.6	191.4	1468	2001	2831	1154	475.6	13.7
$d_0sig$	1338	30.29	235.1	244.3	131.4	854.9	1928	1786	1044	248.1	15.1
$z_0sig$	1287	19.54	105.0	234.7	81.77	408.2	1845	909.9	883.4	244.6	16.6
$M_{\tau^+\tau^-}$	1286	19.39	103.2	234.7	72.05	349.1	1837	883.5	883.4	243.9	16.7
$E_{\tau^+\tau^-}$	1282	19.39	103.0	234.7	72.05	324.7	1836	873.2	883.4	243.9	16.7
$M_{colapp}$	1065	3.074	18.76	47.94	10.28	72.83	616.9	150.8	137.0	0.746	23.1
$E_{colapp}$	1062	2.454	18.01	46.72	10.28	71.27	612.1	93.05	93.52	0.454	23.7
$M_{recoil}$	1026	2.144	14.54	21.24	9.938	57.07	366.3	39.64	43.31	0.161	25.8

## 5 Summary

We evaluated the measurement accuracy of the branching ratio of the  $H \rightarrow \tau^+\tau^-$  mode at  $\sqrt{s} = 250$  GeV at the ILC with ILD\_00 detector model. We assumed  $M_H = 120$  GeV,  $\text{Br}(H \rightarrow \tau^+\tau^-) = 8.0\%$ ,  $\int L dt = 250 \text{ fb}^{-1}$ , and the polarization  $P(e^+, e^-) = (+0.3, -0.8)$ . The obtained values were summarized in Table 4.

Table 4: The analysis results of  $\sqrt{s} = 250$  GeV.

mode	$Z \rightarrow e^+e^-$	$Z \rightarrow \mu^+\mu^-$	$Z \rightarrow q\bar{q}$
significance	$8.0\sigma$	$8.8\sigma$	$25.8\sigma$

From these results, the combined significance was calculated to be  $\sqrt{8.0^2 + 8.8^2 + 25.8^2} = 28.4\sigma$ . Therefore, the measurement accuracy  $\Delta(\sigma \cdot \text{Br})/(\sigma \cdot \text{Br})$  was calculated to be  $\Delta(\sigma \cdot \text{Br})/(\sigma \cdot \text{Br}) = 1/28.4 = 3.5\%$ .

The results are extrapolated to the case of  $M_H = 125$  GeV by scaling the signal yields by the  $e^+e^- \rightarrow ZH$  cross section and the branching ratio  $\text{Br}(H \rightarrow \tau^+\tau^-) \rightarrow 6.32\%$  [9]. We assumed that the selection efficiencies the same. The results are summarized in Table 5.

Table 5: The results of the extrapolation to  $M_H = 125$  GeV.

$Z \rightarrow e^+e^-$	$Z \rightarrow \mu^+\mu^-$	$Z \rightarrow q\bar{q}$	Combined	$\frac{\Delta(\sigma \cdot \text{Br})}{\sigma \cdot \text{Br}}$
$6.8\sigma$	$7.4\sigma$	$21.9\sigma$	$24.1\sigma$	4.2 %







## B Histograms of cut variables

### B.1 $Z \rightarrow e^+e^-$ mode

Figures 17 - 28 show the histograms of cut variables. The blue lines in all histograms show the signal process  $ZH \rightarrow e^+e^-\tau^+\tau^-$ .

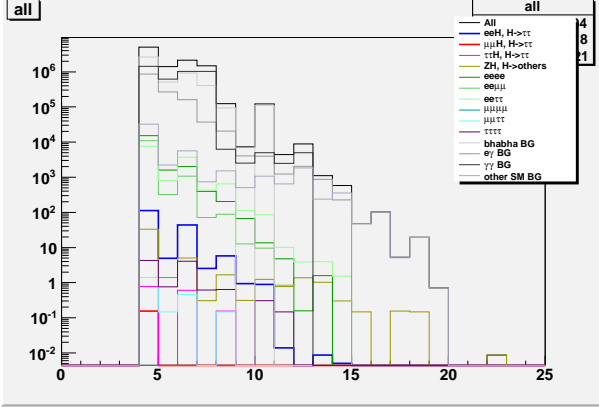


Figure 17: Number of tracks  $\leq 8$ .

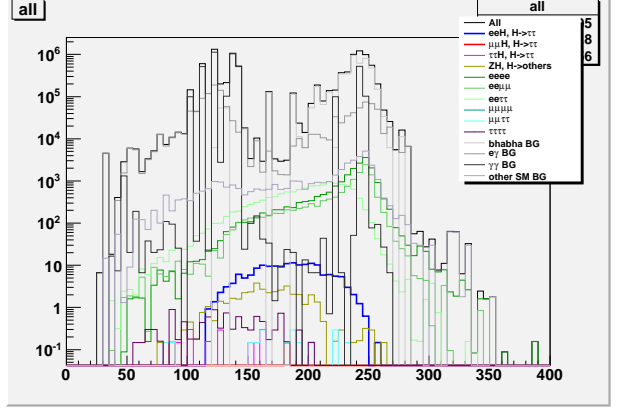


Figure 18:  $115 \text{ GeV} < E_{\text{vis}} < 230 \text{ GeV}$ .

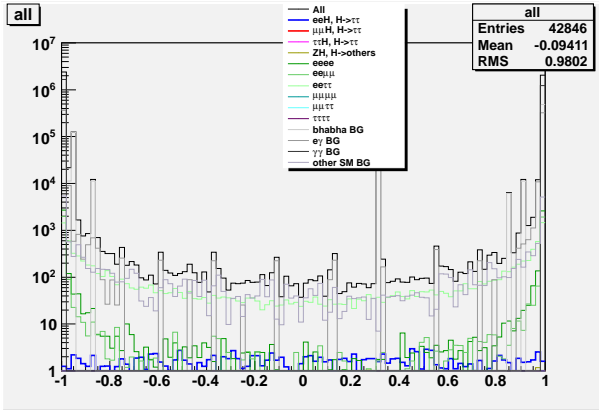


Figure 19:  $|\cos \theta_{\text{miss}}| < 0.99$ .

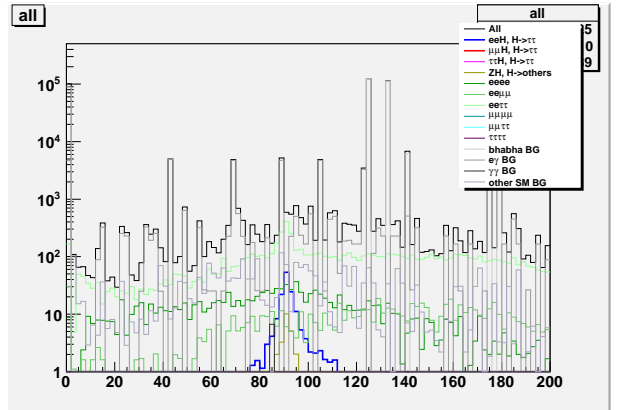


Figure 20:  $81 \text{ GeV} < M_Z < 113 \text{ GeV}$ .

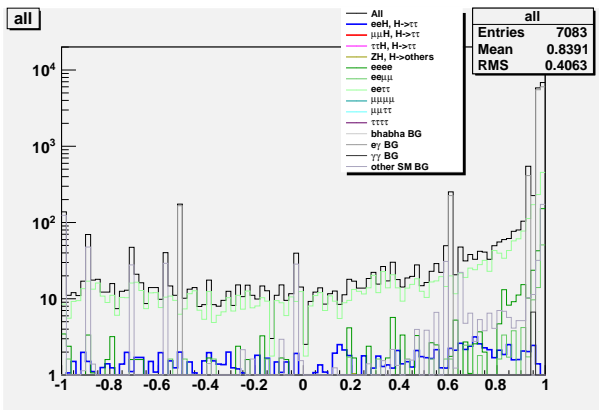


Figure 21:  $\cos \theta_{e^-} < 0.92$ .

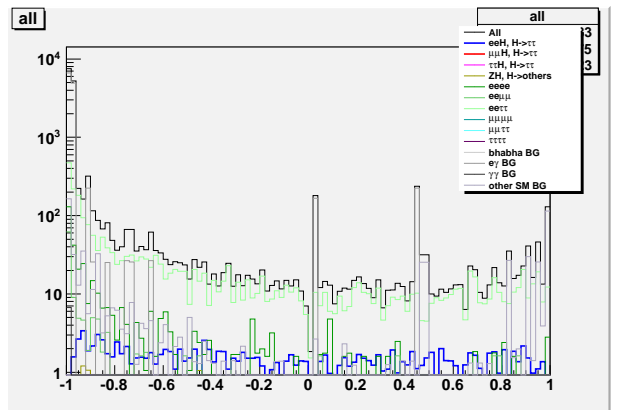


Figure 22:  $\cos \theta_{e^+} > -0.92$ .

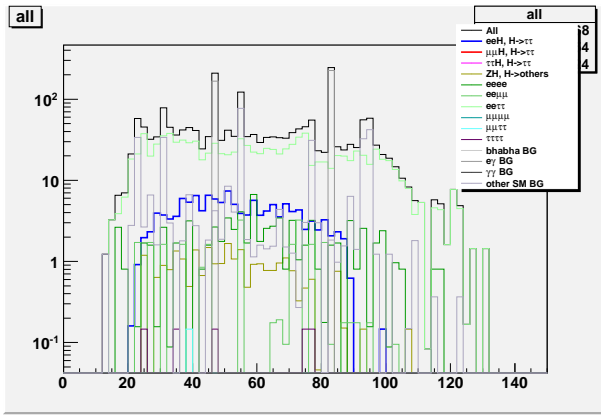


Figure 23:  $E_{e^-} < 90$  GeV.

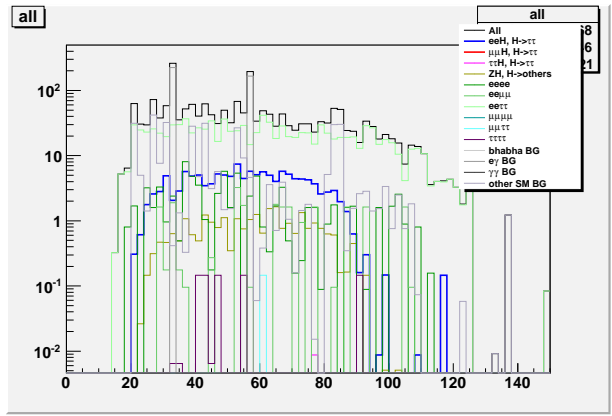


Figure 24:  $E_{e^+} < 90$  GeV.

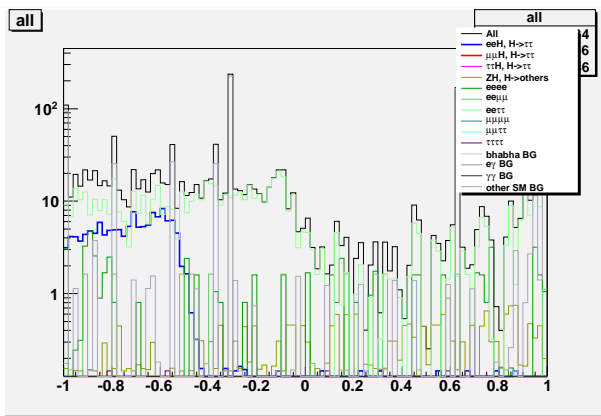


Figure 25:  $\cos \theta_{\tau^+\tau^-} < -0.45$ .

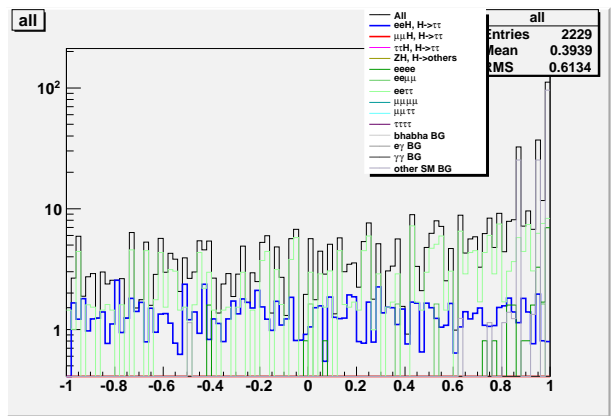


Figure 26:  $\cos \theta_{\tau^-} < 0.92$ .

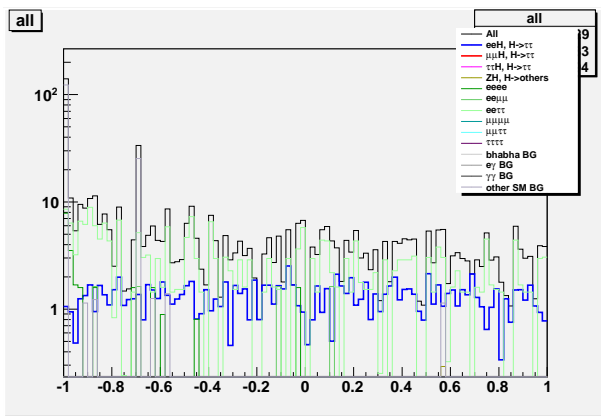


Figure 27:  $\cos \theta_{\tau^+} > -0.92$ .

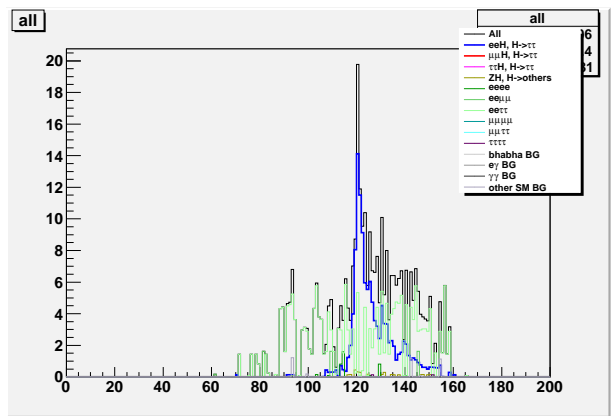


Figure 28:  $116 \text{ GeV} < M_{\text{recoil}} < 142 \text{ GeV}$ .

## B.2 $Z \rightarrow \mu^+\mu^-$ mode

Figures 29 - 36 show the histograms of cut variables. The red lines in all histograms show the signal process  $ZH \rightarrow \mu^+\mu^-\tau^+\tau^-$ .

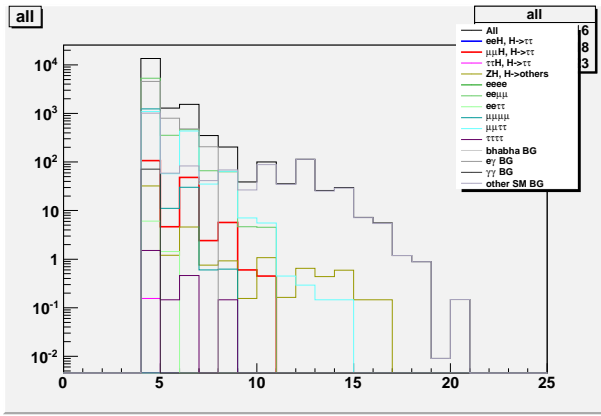


Figure 29: Number of tracks  $\leq 8$ .

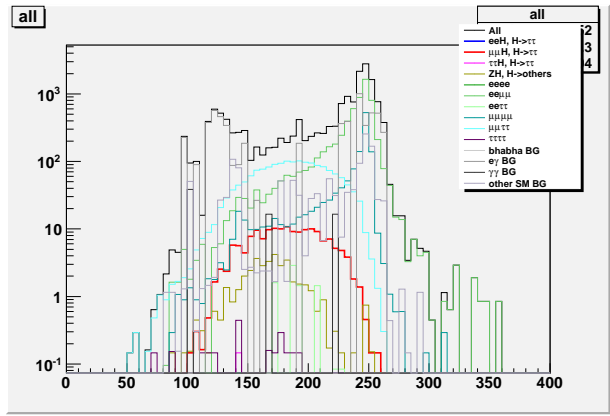


Figure 30:  $115 \text{ GeV} < E_{\text{vis}} < 235 \text{ GeV}$ .

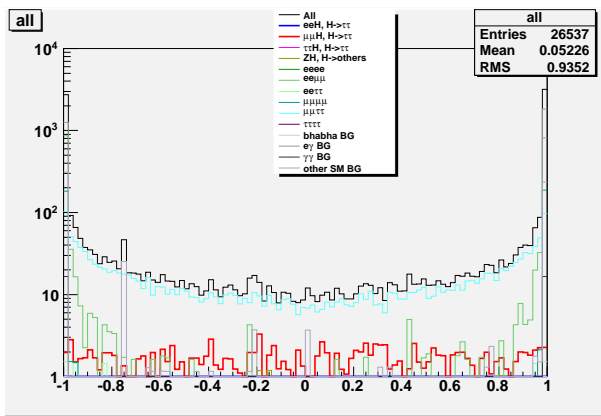


Figure 31:  $|\cos \theta_{\text{miss}}| < 0.98$ .

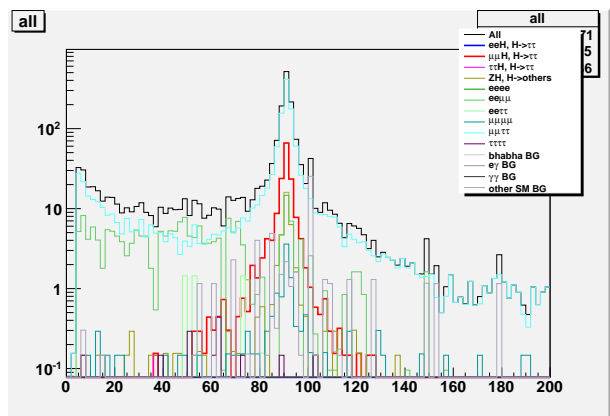


Figure 32:  $72 \text{ GeV} < M_Z < 107 \text{ GeV}$ .

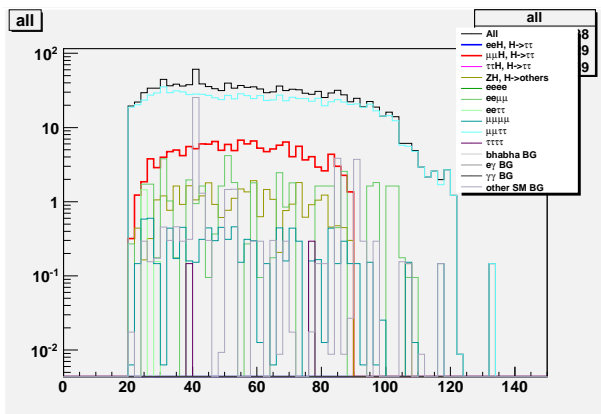


Figure 33:  $E_{e^-} < 90 \text{ GeV}$ .

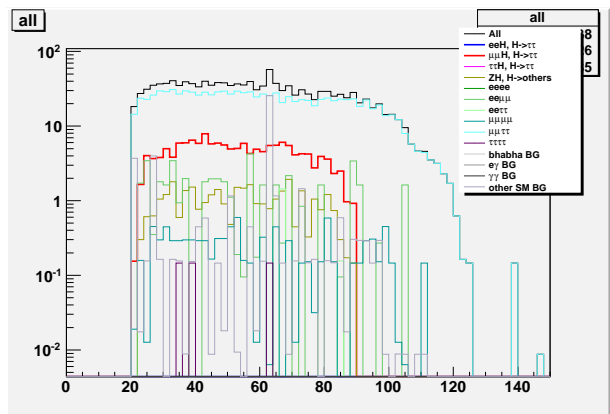


Figure 34:  $E_{e^+} < 90 \text{ GeV}$ .

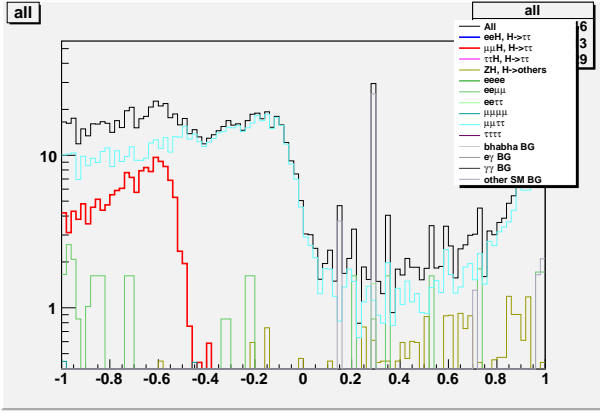


Figure 35:  $\cos \theta_{\tau^+\tau^-} < -0.5$ .

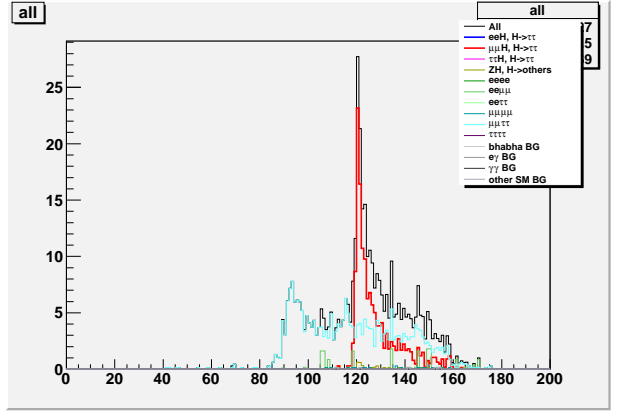


Figure 36:  $118 \text{ GeV} < M_{\text{recoil}} < 143 \text{ GeV}$ .

### B.3 $Z \rightarrow q\bar{q}$ mode

Figures 37 - 49 show the histograms of cut variables. The blue lines in all histograms show the signal process  $ZH \rightarrow q\bar{q}\tau^+\tau^-$ .

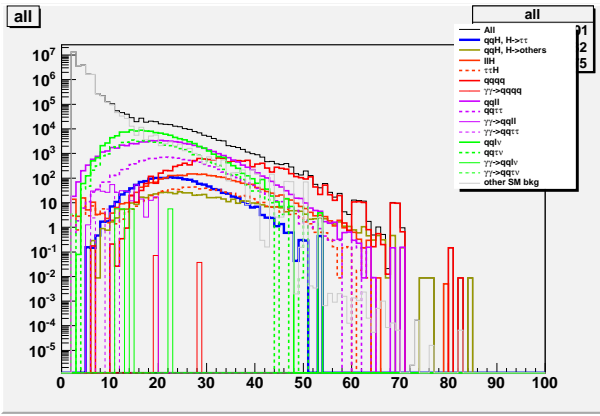


Figure 37:  $9 \leq \text{number of tracks} < 50$ .

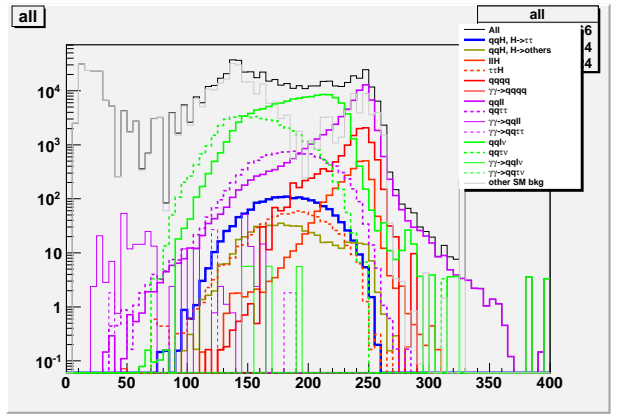


Figure 38:  $110 \text{ GeV} < E_{\text{vis}} < 235 \text{ GeV}$ .

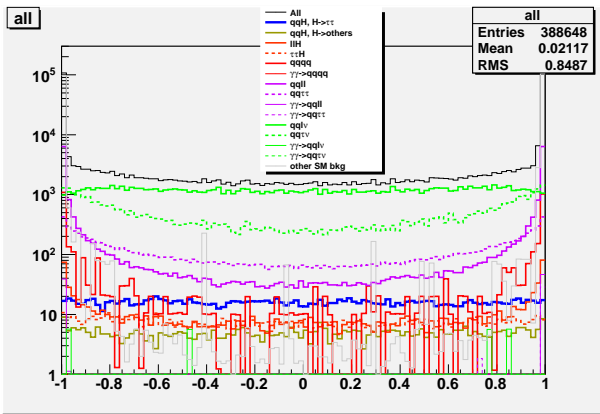


Figure 39:  $|\cos \theta_{\text{miss}}| < 0.98$ .

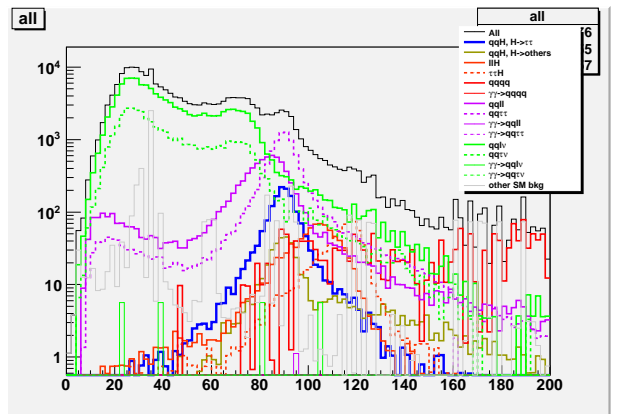


Figure 40:  $77 \text{ GeV} < M_Z < 135 \text{ GeV}$ .

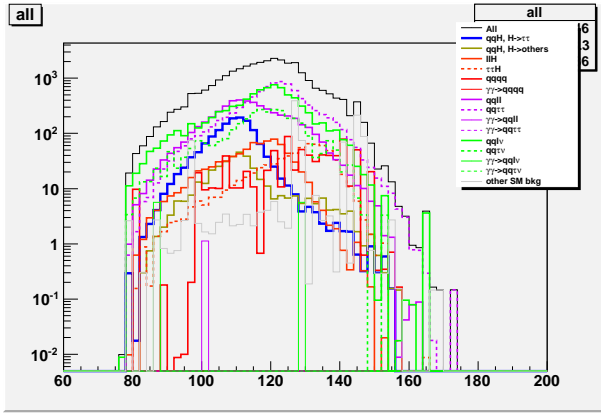


Figure 41:  $80 \text{ GeV} < E_Z < 135 \text{ GeV}$ .

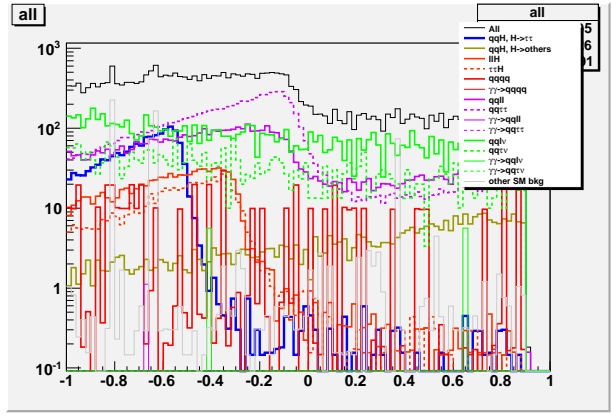


Figure 42:  $\cos \theta_{\tau^+\tau^-} < -0.5$ .

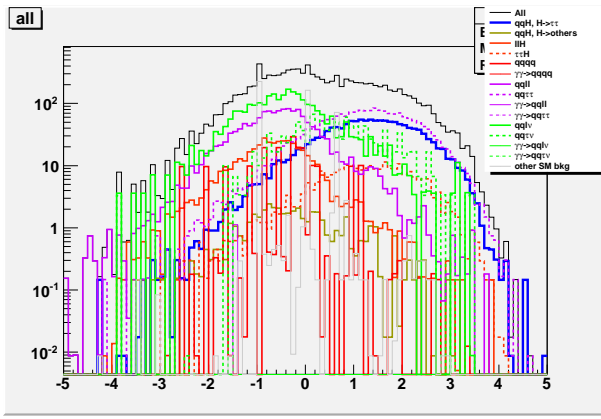


Figure 43:  $\log_{10}(|d_0/\sigma(d_0)|)(\tau^-) + \log_{10}(|d_0/\sigma(d_0)|)(\tau^+) > -0.7$ .

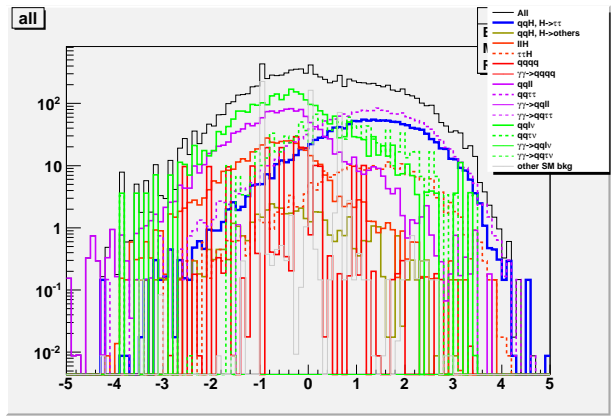


Figure 44:  $\log_{10}(|z_0/\sigma(z_0)|)(\tau^-) + \log_{10}(|z_0/\sigma(z_0)|)(\tau^+) > -0.1$ .

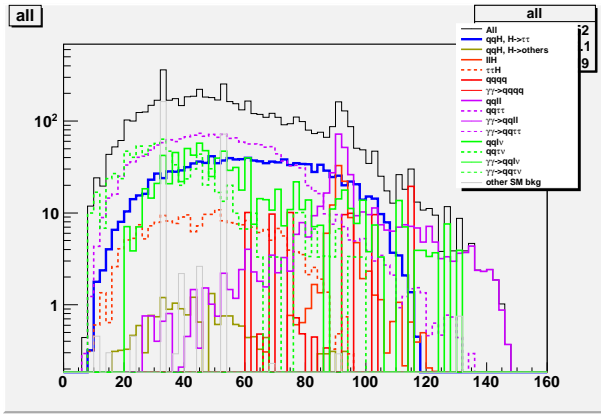


Figure 45:  $M_{\tau^+\tau^-} < 115 \text{ GeV}$ .

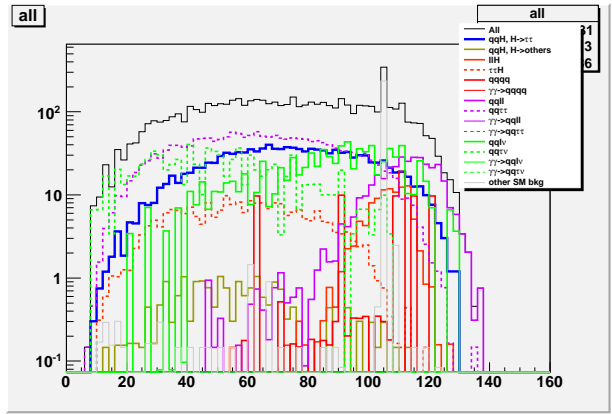


Figure 46:  $E_{\tau^+\tau^-} < 125 \text{ GeV}$ .

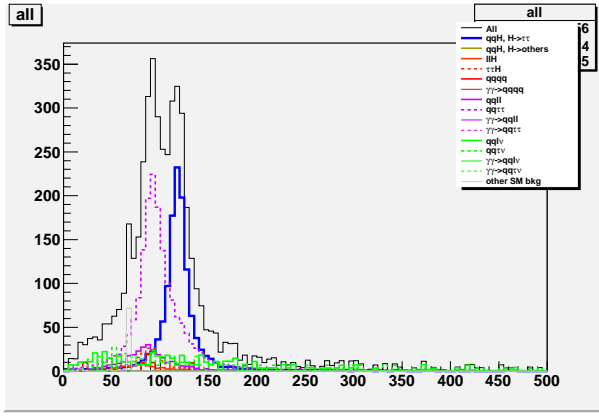


Figure 47:  $100 \text{ GeV} < M_{\text{colapp}} < 170 \text{ GeV}$ .

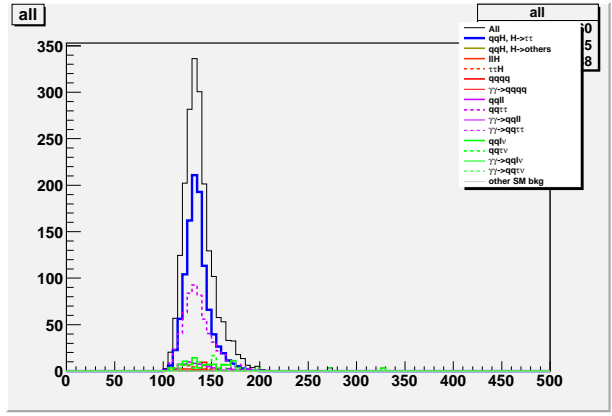


Figure 48:  $100 \text{ GeV} < E_{\text{colapp}} < 280 \text{ GeV}$ .

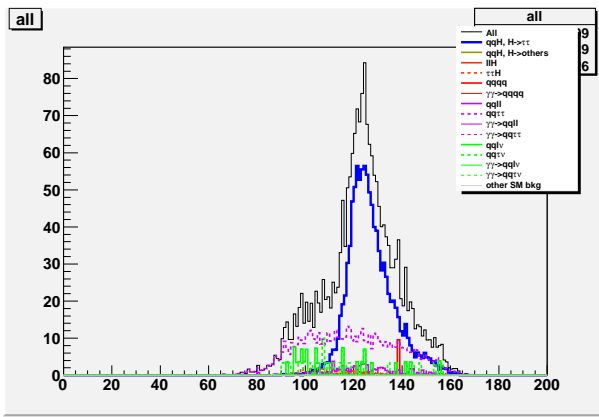


Figure 49:  $112 \text{ GeV} < M_{\text{recoil}} < 160 \text{ GeV}$ .

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