Higgs Branching Fraction Study in ILC

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Precise measurement of the Higgs boson properties are important issues for the International Linear Collider (ILC) project to understand the particles mass generation mechanism which strongly related to the coupling with the Higgs boson. Large Hadron Collider (LHC) [1] experiments exclude the large area of the predicted Higgs mass region and their results indicate that Higgs boson mass will be light. Even if LHC discovers the Higgs like particle by the end of 2012, Higgs will be identified by the high precision measurement of the Higgs boson properties in ILC and also Higgs measurement verifies the correctness of standard model (SM) or gives some hints toward its beyond. In this study, we evaluate the measurement accuracies of Higgs branching fraction to the $H \rightarrow b\bar{b}$, $c\bar{c}$ and gg at the center-of-mass energy of 250 and 350 GeV.

1 Introduction

Higgs branching ratio (BR) measurement is one of the important issue of the International Linear Collider (ILC) project, which is strongly related to the coupling strength with particles and reveal their mass generation mechanism. Even the Large Hadron Collider (LHC) [1] will discover the Higgs boson in a few year, ILC can confirm whether that is the standard model (SM) predicted one or not, and find some hints toward its beyond. LHC experiment accumulate the total integrated luminosity up to 5 fb⁻¹ by the end of 2011, LHC gradually exclude the heavy Higgs mass region and indicates the light Higgs from

Table 1: Expected SM Higgs BR at the Higgs mass of 120 GeV in PYTHIA [2].

| H decay mode | BRs at $M_H = 120 \text{ GeV}$ |
|-----------------------|----------------------------------|
| $H \rightarrow bb$ | 65.7% |
| $H \to WW^*$ | 15.0% |
| $H \to \tau^+ \tau^-$ | 7.9% |
| $H \rightarrow gg$ | 5.5% |
| $H \to c\bar{c}$ | 3.6% |

the combined results of ATLAS [3] and CMS [4] ($115 \leq M_H \leq 135$ GeV). In this region, we obtain the maximum production cross section around the center-of-mass energy (\sqrt{s}) of 250 GeV and Higgs boson BR significantly varies depending on the Higgs mass, especially main decay channel shifts from $H \to b\bar{b}$ to WW^* around the $M_H = 140$ GeV, as shown in Fig. 1 (a) and (b). In this study, we evaluate the measurement accuracies of the Higgs BRs of $H \to b\bar{b}$, $c\bar{c}$ and gg with assuming the Higgs mass of 120 GeV and integrated luminosity (\mathcal{L}) of 250 fb⁻¹, using the International Large Detector (ILD) [5] full simulation at $\sqrt{s} = 250$ GeV. In addition, we also consider the operation at the CM energy of 350 GeV, which increase the contribution of W/Z fusion process but Z/H will be boosted, for taking into account the energy staging option in ILC project. Considering the LHC results, we also evaluate the Higgs BR accuracy at the indicated mass region by extrapolating the $M_H = 120$ GeV result.

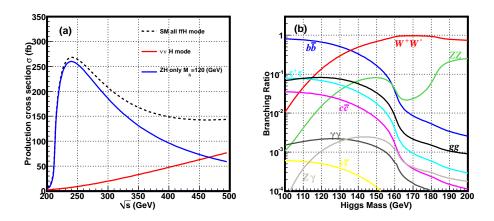


Figure 1: (a) Higgs production cross section as a function of CM energy and (b) SM predicted Higgs BR as a function of Higgs mass.

2 Signal and Backgrounds

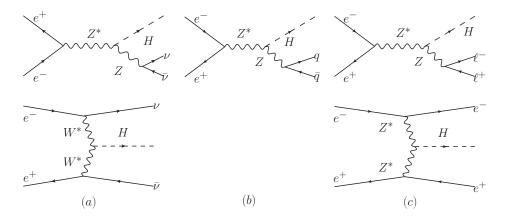


Figure 2: Higgs production processes categorized with Z decay channels: (a) neutrino $(\nu \bar{\nu} H)$, (b) hadronic $(q\bar{q}H)$ and (c) leptonic $(\ell^+\ell^-H)$.

Figure 2 shows the Higgs production diagrams and ZH analysis procedures are categorized with the Z decay channels; $ZH \rightarrow \nu\bar{\nu}H$ (Neutrino), $q\bar{q}H$ (Hadronic) and $\ell^+\ell^-H$ (Leptonic) which mainly form di-jet, four-jet and di-lepton+di-jet final states, respectively. We also consider the W/Z-fusion process for the neutrino and leptonic channels which has larger contribution at $\sqrt{s} = 350$ GeV. Assuming the $M_H = 120$ GeV, we obtain the maximum production cross section around the $\sqrt{s} = 250$ GeV through the Higgsstrahlung (ZH) process and mainly decays to b quarks ($H \rightarrow b\bar{b}$). In order to maximize the Higgs production cross section, we employ the left-handed electron beam polarization; $(e^+, e^-) = (+30\%, -80\%)$, at the CM energies of both 250 and 350 GeV with the integrated luminosity of 250 fb⁻¹. As background, we consider the following 2f and 4f final state SM Backgrounds: $e^+e^- \rightarrow W^+W^-$, ZZ and $q\bar{q}$. In addition, we also take into account for the $e^+e^- \rightarrow t\bar{t}$ background only for the $\sqrt{s} = 350$ GeV.

3 Analysis Framework

Since ZH final state forms multi-jet, thus jet clustering, jet energy resolution and quark flavor-tagging are crucial for the Higgs hadronic decay channels $(H \rightarrow bb, c\bar{c} \text{ and } qq)$ analysis. In order to achieve the best jet energy resolution, ILD adopt the Particle Flow Algorithm (PFA) [6], which can achieve the best jet energy resolution; charged tracks energy is measured by tracker instead of the calorimeter and only neutral particles energy is measured by calorimeter with avoiding the cluster overlapping and double counting. ILD detector design is well-suitable for the best PFA performance with adopting the finely segmented calorimeter, large tracker radius and strong magnetic field. For the simulation study, we use the ilcsoft v01_06 [7] ILC common software package. At fist we generate the MC event samples with the Whizard [8]. Then we perform the ILD full detector simulation with the Mokka [9] package assuming the ILD detector model (ILD_00). Generated hits are digitized and reconstructed with Marlin [10] package and perform the PFA (PandraPFA [6]). We also employ the LCFIVTX [11] flavor tagging package embedded in Marlin to identify the quark flavor of $H \to bb$ and $c\bar{c}$. For the mass production of the full simulation samples, we use the GRID [12] resources for the ILD LOI study [5] and finally saved as ILC common file format (LCIO [13]).

4 Event Reconstruction and Background Reduction

4.1 Neutrino Channel $(\nu \bar{\nu} H)$

4.1.1 Jet Reconstruction

For the $\nu \bar{\nu} H$ channel analysis, we apply the jet reconstruction forcibly merged into the two-jet which comes from $H \to b\bar{b}$, $c\bar{c}$ and gg.

4.1.2 Background Reduction

In this channel, $ZZ \rightarrow \nu\nu qq$ and $WW \rightarrow \nu\ell qq$ will be the main backgrounds. To suppress these backgrounds, at first, we apply the missing mass cut; $80 < M_{miss} < 140$ GeV to suppress the ZZ to leptonic or hadronic decay backgrounds, since M_{miss} should be consistent with the Z mass in $Z \rightarrow \nu\bar{\nu}$ channel. Then we use the following kinematic variables cut: transverse momentum; $20 < P_t < 70$ GeV, longitudinal momentum; $|P_\ell| < 60$ GeV and maximum momentum; $P_{max} < 30$ GeV to suppress the $q\bar{q}$ background. Number of charged tracks cut; $N_{chd} > 10$ well reduce the backgrounds including energetic leptons. We also apply the jet clustering y-value cuts; $0.2 < Y_{12} < 0.8$ and $Y_{23} < 0.02$, which are the y-value thresholds from two- to one-jet or two- to three-jets and different of number of jets backgrounds are reduced with these cut. Finally we apply the di-jet mass cut; $80 < M_{jj} < 130$ GeV which correspond to Higgs mass. After applying all the cuts, we apply the likelihood variable cut using following input variables; M_{miss} , number of reconstructed particles (N_{PFO}) , P_{max} , P_l and M_{jj} , and we select the LR > 0.375 where we obtain the maximum signal significance. In order to optimize the cut positions for the analysis at $\sqrt{s} = 350$ GeV, we change the following variables cut positions: $50 < M_{miss} < 240$ GeV, $10 < P_t < 140$ GeV, $|P_\ell| < 130$ GeV and LR > 0.15.

Table 2: Summary of the background reduction in $\nu \bar{\nu} H$ channel at the $\sqrt{s} = 250$ and 350 GeV.

| $\sqrt{s} \; (\text{GeV})$ | | 250 | 350 | | |
|----------------------------|-------|----------|-------------|----------|--|
| Cut | Sig. | Bkg. | Sig. | Bkg. | |
| Gen. | 19360 | 44827100 | 26307 | 20855900 | |
| All cuts | 6731 | 19058 | 12338 | 71918 | |
| LR cut | 4753 | 3593 | 9302 | 10029 | |
| Significance (Eff.) | 52.0 | (24.5%) | 66.9(35.4%) | | |

4.2 Hadronic Channel $(q\bar{q}H)$

4.2.1 Jet Reconstruction and Pairing

For $q\bar{q}H$ channel, we apply the four-jet reconstruction forcibly. Then we calculate following χ^2 value to determine the Z or H jet pair candidates:

$$\chi^2 = \left(\frac{M_{12} - M_H}{\sigma_H}\right)^2 + \left(\frac{M_{34} - M_Z}{\sigma_Z}\right),\tag{1}$$

where M_{12} and M_{34} are reconstructed di-jet invariant masses and $\sigma_{Z/H}$ are the width of Z and H mass distribution. We select the minimum χ^2 jet pairs as the best candidate of Z and H.

4.2.2 Background Reduction

After the jet clustering and pairing, we apply the background reduction. At first we apply the χ^2 cut to reduce the wrong combination pairs; $\chi^2 < 10$. Then we require the following cuts to suppress the leptonic events: the number of charged tracks cut; $N_{chd} > 4$ and y-value cut Y_{34} cut, which is a y-value threshold from three- to four-jet; $-\log Y_{34} < 2.7$. As event shape cuts, we employ the following variables cut: thrust and its cosine of thrust angle; thrust < 0.9, $|\cos \theta_{thrust}| < 0.9$ and the angle between the Higgs candidate jets; $105 < \theta_H < 160^\circ$ to suppress the $q\bar{q}$ backgrounds. Finally we apply the di-jet mass cut M_{12} and M_{34} and likelihood variable cut which is calculated with following input variables; thrust, M_Z , M_H , θ_H . We select the likelihood variable cut as LR > 0.2. For the $\sqrt{s} = 350$ GeV, we optimize the cut position as follows; thrust < 0.85, $70 < \theta_H < 120^\circ$, $80 < M_Z < 100$ GeV, $105 < M_H < 130$ GeV and LR > 0.1.

4.3 Leptonic Channel $(\ell^+\ell^-H)$

For the $\ell^+\ell^- H$ channel analysis, at first we identify the di-lepton, then we apply the di-jet reconstruction forcibly for remaining particles.

| $\sqrt{s} \; (\text{GeV})$ | | 250 | 350 | | |
|----------------------------|-----------|----------|-------|----------|--|
| Cut | Sig. Bkg. | | Sig. | Bkg. | |
| Gen. | 52507 | 44827100 | 36099 | 21222700 | |
| All cuts | 16350 | 411785 | 9447 | 44400 | |
| LR cut | 13726 | 166807 | 8686 | 25393 | |
| Significance (Eff.) | 32.3 | (26.1%) | 47.1 | (24.1%) | |

Table 3: Summary of the background reduction in $q\bar{q}H$ channel at the $\sqrt{s} = 250$ and 350 GeV.

4.3.1 Di-lepton Identification

We apply the following di-lepton identification for electrons and muons from the different aspects in the energy deposition in the calorimeter; $E_{ECAL}/E_{Total} > 0.9$ and $0.7 < E_{Total}/P < 1.2$ for electrons and $E_{ECAL}/E_{Total} < 0.5$ and $E_{Total}/P < 0.4$ for muons, where E_{ECAL} , E_{Total} and P denote the ECAL energy associated with a track, total energy deposit in whole calorimeter and track momentum.

4.3.2 Background Reduction

After the di-lepton identification, we apply the background reduction with following cut variables, which is summarized on the Table 4. At first we apply the di-lepton mass $(M_{\ell\ell})$ cut which should be consistent with the Z mass: $70 < M_{\ell\ell} < 110$ GeV for electron and $70 < M_{\ell\ell} < 100$ GeV for muons, respectively. Then we apply the Z flight direction cut: $|\cos \theta_Z| < 0.8$ to suppress the forward region backgrounds. Finally we require the di-jet mass (M_{jj}) and recoil mass (M_{rec}) cuts to select the Higgs candidate signal: $100 < M_{jj} < 140$ GeV, $70 < M_{rec} < 140$ GeV for electron; $115 < M_{jj} < 140$ GeV and $70 < M_{rec} < 140$ GeV for muon, respectively.

| $\sqrt{s} \; (\text{GeV})$ | 250 350 | | | 350 | |
|-----------------------------|--------------|---------|--------------|---------|--|
| Cut | Sig. | Bkg. | Sig. | Bkg. | |
| Gen. (e) | 3137 | 4512520 | 2740 | 3822410 | |
| Gen. (μ) | 2917 | 4512520 | 1789 | 3822410 | |
| All cuts (e) | 1184 | 1607 | 567 | 590 | |
| All cuts (μ) | 1365 | 983 | 638 | 465 | |
| Significance (Eff.) (e) | 22.4 (37.8%) | | 16.7 | (20.7%) | |
| Significance (Eff.) (μ) | 28.2 | (46.8%) | 19.2 (35.7%) | | |

Table 4: Summary of the background reduction in $\ell^+\ell^-H$ channel at the $\sqrt{s} = 250$ and 350 GeV.

5 Measurement Accuracy of Branching Fraction

5.1 Template Fitting

In order to evaluate the measurement accuracies of the Higgs branching fraction (BR), we employ the template fitting method [14]. At first, we prepare the flavor-likeness template samples which is calculated from the LCFIVTX output $x_{1,2}$:

$$x - likeness \equiv \frac{x_1 x_2}{(1 - x_1)(1 - x_2)},$$
(2)

where $x_{1,2}$ represents the flavor tagging output from LCFIVTX for di-jet. We assume the Poisson statistics (P_{ink}) for each bin (i, j, k) of the template samples:

$$P_{ijk} = \frac{X^{\mu} e^{-\mu}}{X!} \quad \left(\mu \equiv \sum N_{ijk}^{template}, \ X \equiv N_{ijk}^{data}\right),\tag{3}$$

where N_{ijk}^{data} is the number of entries in (i, j, k) bin. $N_{ijk}^{template}$ represents the sum of the number of entries at bin (i, j, k) in each template sample:

$$N_{ijk}^{template} = \sum_{s=bb,cc,gg,bkg} r_s \cdot N_{ijk}^s,\tag{4}$$

where N_{ijk}^s represents the number of entries at the (i, j, k) bin in $H \to b\bar{b}$, $c\bar{c}$, gg and background template sample (N_{ijk}^{bkg}) , which includes the number of entries of SM background and Higgs to none hadronic decays. r_s represents the fitting parameters of r_{bb} , r_{cc} , r_{gg} and r_{bkg} , where they are the ratios of number of entries in Higgs hadronic decays of $H \to b\bar{b}$, $c\bar{c}$ and gg after the background reduction to the entries predicted from the SM Higgs BR. r_{bkg} is a normalization factor for the SM background and other Higgs none hadronic decays, which is fixed to be 1 from the assumption that the SM backgrounds are well understood. Finally we apply the template fitting with minimizing the following log-likelihood variable L calculated from the product of the probability P_{ijk} in each bin:

$$L = -\log\left(\prod_{i,j,k} P_{ijk}\right) = -\sum_{i,j,k} (\log P_{ijk}).$$
(5)

5.2 Measurement accuracies of BR

To evaluate the measurement accuracies of the sigma times BRs for bb, cc and gg, we apply the 1000 times template fitting Toy-MC and obtain the fitted results of r_s (s = bb, cc, gg):

$$\sigma_{ZH} \cdot BR(H \to s) = r_s \times \sigma_{ZH}{}^{SM} \cdot BR(H \to s){}^{SM}, \tag{6}$$

where σ_{ZH} is a Higgs production cross section, σ_{ZH}^{SM} and $BR(H \to s)^{SM}$ are cross section and BR of $H \to s$ predicted in SM. From the Eq. 6, we obtain the measurement accuracies of Higgs BR from following equation:

$$\frac{\Delta\sigma BR(H\to s)}{\sigma BR} = \sqrt{\left(\frac{\Delta r_s}{r_s}\right)^2 + \left(\frac{\Delta\sigma_{ZH}}{\sigma_{ZH}}\right)^2},$$

here we assume the 2.5% of cross section measurement uncertainty $(\Delta \sigma_{ZH}/\sigma_{ZH})$ estimated from the recoil mass study [15]. Summary tables of the measurement accuracies of Higgs BR are shown in Table 5 and 6 for $\sqrt{s} = 250$ and 350 GeV, respectively.

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| | $\nu \bar{\nu} H$ | $q\bar{q}H$ | e^+e^-H | $\mu^+\mu^-H$ | comb. |
|--------------------------------------|-------------------|-------------------|------------------|-----------------|-------------------|
| r_{bb} | $1.00{\pm}0.016$ | $1.00{\pm}0.015$ | 1.00 ± 0.039 | 1.00 ± 0.33 | $1.00{\pm}0.012$ |
| r_{cc} | $1.00{\pm}0.12$ | 1.00 ± 0.12 | 0.98 ± 0.29 | 1.01 ± 0.24 | $1.00 {\pm} 0.09$ |
| r_{gg} | $0.99 {\pm} 0.14$ | $1.00 {\pm} 0.13$ | 0.99 ± 0.35 | 1.00 ± 0.21 | $1.00{\pm}0.10$ |
| $\sigma BR(bb)/\sigma^{SM}$ (%) | 65.7 ± 1.1 | 65.7 ± 1.0 | 65.7 ± 2.6 | 65.7 ± 2.2 | $65.7 {\pm} 0.7$ |
| $\sigma BR(cc)/\sigma^{SM}$ (%) | $3.59 {\pm} 0.43$ | $3.61 {\pm} 0.44$ | 3.53 ± 1.03 | 3.63 ± 0.85 | $3.60 {\pm} 0.31$ |
| $\sigma BR(gg)/\sigma^{SM}$ (%) | $5.46 {\pm} 0.76$ | $5.48 {\pm} 0.76$ | 5.45 ± 1.94 | 5.49 ± 1.14 | 5.47 ± 0.54 |
| $\Delta\sigma BR(bb)/\sigma BR~(\%)$ | 3.0 | 2.9 | 4.7 | 3.3 | 2.7 |
| $\Delta\sigma BR(cc)/\sigma BR~(\%)$ | 12.2 | 12.3 | 29.3 | 23.5 | 8.9 |
| $\Delta\sigma BR(gg)/\sigma BR~(\%)$ | 14.2 | 14.1 | 35.6 | 20.7 | 10.2 |

Table 5: Summary of template fitting results and BR measurement accuracies at the $\sqrt{s}=250$ GeV.

Table 6: Summary of template fitting results and BR measurement accuracies at the $\sqrt{s} = 350$ GeV.

| | $\nu \bar{\nu} H$ | $q\bar{q}H$ | e^+e^-H | $\mu^+\mu^-H$ | comb. |
|--------------------------------------|-------------------|-------------------|------------------|------------------|--------------------|
| r_{bb} | $1.00{\pm}0.012$ | $1.00{\pm}0.015$ | 1.00 ± 0.056 | 1.00 ± 0.051 | $1.00 {\pm} 0.010$ |
| r_{cc} | $1.00{\pm}0.10$ | $0.99 {\pm} 0.11$ | 1.02 ± 0.26 | 1.02 ± 0.32 | $1.00 {\pm} 0.07$ |
| r_{gg} | $1.00 {\pm} 0.10$ | $1.00 {\pm} 0.13$ | 0.97 ± 0.35 | 0.97 ± 0.35 | $1.00 {\pm} 0.08$ |
| $\sigma BR(bb)/\sigma^{SM}$ (%) | $65.7 {\pm} 0.8$ | 65.7 ± 1.0 | 65.6 ± 3.68 | 65.6 ± 3.32 | 65.7 ± 0.6 |
| $\sigma BR(cc)/\sigma^{SM}$ (%) | $3.60 {\pm} 0.35$ | $3.68 {\pm} 0.26$ | 3.68 ± 0.94 | 3.66 ± 1.16 | $3.59 {\pm} 0.26$ |
| $\sigma BR(gg)/\sigma^{SM}$ (%) | $5.48 {\pm} 0.53$ | 5.49 ± 0.74 | 5.32 ± 1.91 | 5.35 ± 1.94 | 5.48 ± 0.43 |
| $\Delta\sigma BR(bb)/\sigma BR~(\%)$ | 2.8 | 2.9 | 6.1 | 5.6 | 2.7 |
| $\Delta\sigma BR(cc)/\sigma BR~(\%)$ | 10.1 | 11.2 | 25.6 | 31.7 | 7.7 |
| $\Delta\sigma BR(gg)/\sigma BR~(\%)$ | 9.9 | 13.7 | 36.0 | 36.3 | 8.2 |

6 Extrapolate to the different Higgs masses

Since latest LHC results indicate the light Higgs $(115 \le M_H \le 140 \text{ GeV})$, we should extend the results at the $M_H = 120 \text{ GeV}$ to the prospective mass region. In order to evaluate the Higgs BR measurement accuracy at other Higgs masses, we extrapolate the results at the $M_H = 120 \text{ GeV}$ with Eq. 7 while we assume the same background reduction efficiency for other mass.

$$\left(\frac{\Delta\sigma BR}{\sigma BR}(s)\right)_{M_H} = \left(\frac{\Delta\sigma BR}{\sigma BR}(s)\right)_{120} \times \sqrt{\frac{\sigma_{120} \cdot BR(s)_{120}}{\sigma_{M_H} \cdot BR(s)_{M_H}}},\tag{7}$$

where σ_{M_H} and $BR(s)_{M_H}$ denote the cross section and BR of $H \to s$ at M_H , as shown in Fig. 3 (a). In addition, we also compile the $H \to WW^* \to 4j$ result [16], even though this study assume the electron right-handed polarization: $(e^+, e^-) = (-30\%, +80\%)$ to suppress the $e^+e^- \to W^+W^-$ backgrounds. Figure 3 (b) shows the expected measurement accuracies of the Higgs BR at each mass and summarized on the Table 7.

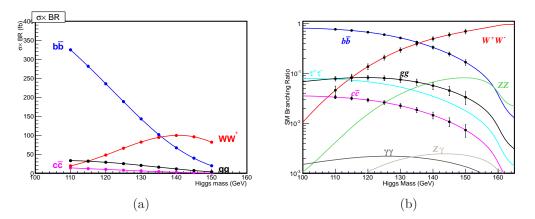


Figure 3: Higgs mass dependence of (a) cross section times Higgs BR and (b) measurement accuracies of the Higgs BR extrapolated from the result at $M_H = 120$ GeV.

7 Conclusion

We evaluate the measurement accuracies of the Higgs BR for $H \rightarrow b\bar{b}$, $c\bar{c}$ and gg channels. With the template fitting analysis, we obtain the measurement accuracies of $\sigma \times BR$ for $H \rightarrow b\bar{b}$, $c\bar{c}$ and gg as 3%, 9% and 10%, respectively. We also estimate the Higgs BR measurement accuracies at the prospective mass region by LHC with extrapolating the results at the Higgs mass of 120 GeV.

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Table 7: Summary of the extrapolated accuracies of the Higgs BR from the result at the $M_H = 120$ GeV, $\sqrt{s} = 250$ GeV to prospective masses calculated by HPROD [17] and HDECAY [18]. Here 2.5% of the σ_{ZH} uncertainty is also included.

| $M_H (\text{GeV})$ | 120 | | | 130 | | | 140 | | |
|--------------------------------|-------|-------------|--------------------------------------|-------|-------------|--------------------------------------|-------|-------------|--------------------------------------|
| $\sigma \; (\mathrm{fb}^{-1})$ | 354.3 | | | 279.9 | | | 203.1 | | |
| Modes | BR(%) | σBR | $\frac{\Delta \sigma BR}{\sigma BR}$ | BR(%) | σBR | $\frac{\Delta \sigma BR}{\sigma BR}$ | BR(%) | σBR | $\frac{\Delta \sigma BR}{\sigma BR}$ |
| $H \to b\bar{b}$ | 66.5 | 235.6 | 2.7% | 51.2 | 143.3 | 3.5% | 33.0 | 67.1 | 5.1% |
| $H \to c\bar{c}$ | 2.9 | 10.4 | 8.1% | 2.3 | 6.3 | 10.4% | 1.5 | 3.0 | 15.2% |
| $H \rightarrow gg$ | 8.2 | 29.2 | 9.0% | 7.5 | 21.0 | 10.6% | 5.7 | 11.5 | 14.3% |
| $H \to W^+ W^-$ | 13.6 | 48.3 | 15.7% | 29.4 | 82.4 | 10.3% | 49.2 | 99.8 | 9.3% |

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8 Bibliography

References

- [1] Large Hadron Collider, http://lhc.web.cern.ch/
- [2] T. Sjöstrand, S. Mrenna and P. Skands, JHEyP05 (2006) 026.
- [3] ATLAS experiment, http://www.atlas.ch/
- [4] CMS experiment, http://cms.web.cern.ch/
- [5] T. Abe et al. [ILD Concept Group Linear Collider Collaboration], "The International Large Detector: Letter of Intent", KEK Report 2009-6
- [6] M. A. Thomson, arXiv:0907.3577v1 [physics.ins-det] (2009)
- [7] http://ilcsoft.desy.de/portal/
- [8] W. Kilian et.al, arXiv: 0708.4233 [hep-ph]
 M. Moretti et.al, arXiv: hep-ph/0102195-rev.
- [9] http://polzope.in2p3.fr:8081/MOKKA/
- [10] http://ilcsoft.desy.de/marlin/
- [11] Nuclear Instruments and Methods in Physics Research Section A, Volume 610, Issue 2, p. 573-589.
- [12] http://lcg.web.cern.ch/lcg/
- $[13]\,$ F. Gade et. al., arXiv:physics/0306114v1 [physics.data-an] (2003)
- $[14]\,$ T. Kuhl and K. Desch, LC-PHSM-2007-001
- $[15]~{\rm H.Li,~arXiv:1007.2999v1}~[{\rm hep-ex}]~(2010)$
- [16] Y. Takubo et. al., arXiv:1011.5805v3 [hep-ph] (2010)
- $\left[17\right]\,$ A. Djouadi, GDR Note GDR-S-012.
- [18] A. Djouadi, arXiv:hep-ph/9704448v1
- [19] http://www-jlc.kek.jp/subg/physics/ilcphys/